

Solvation of Alkali Metal Cations of Arylmethyl Anions by Ethereal Solvents

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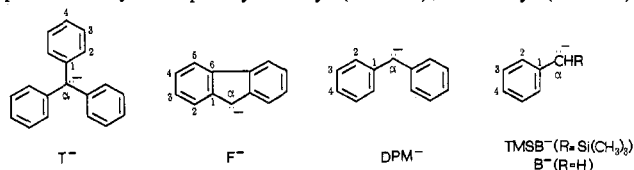
Abstract: The ion pair equilibria for triphenylmethyl, fluorenyl, diphenylmethyl, α -trimethylsilylbenzyl, and benzyl alkali metal anions dissolved in 2-methyltetrahydrofuran, tetrahydrofuran, and dimethoxyethane have been investigated with variable temperature carbon-13 magnetic resonance. The α -carbon chemical shifts for contact ions are very sensitive to cation and are proportional to the reciprocal of the interionic distance. Large α -carbon chemical shift changes are observed for the lithium and sodium salts when the change from contact ions at high temperatures to solvent-separated ions at low temperatures is within the experimentally accessible temperature range. The α -carbon chemical shifts for lithium and sodium solvent-separated ions are cation independent and comparison of these shifts to the contact shifts of the larger alkali metal carbanions gives an estimate of the interionic distance in the solvent-separated state. Little or no temperature-dependent chemical shift change is observed for the potassium, rubidium, or cesium salts either because they are contact ions at all temperatures or because carbon-13 is insensitive to solvation if it occurs. Entropies and enthalpies of solvation were determined for those salts that undergo approximately 40% of the change from contact to solvent-separated ions within the accessible temperature range.

Introduction

Ultraviolet absorption spectroscopy has been used extensively to characterize the equilibria between contact and solvent-separated ion pairs.^{1,2} Fluorenyl salts have been used to determine the thermodynamics of solvation and the effects of cation and solvent on the equilibria.³⁻⁵ Relatively few other carbanions have been studied because of overlap of the ultraviolet maxima (diphenylalkylmethyl⁶ and 1,3-diphenylallyl^{7,8}). Proton magnetic resonance has also been used to qualitatively study ion-pair equilibria.^{4,7,9-16} Shift changes *indirectly* reflect π -electron density changes and ¹H NMR can be used to study a wider variety of carbanion systems (fluorenyl,^{4,9,10-17} indenyl,^{11,12,17} cyclopentadienyl,¹⁷ triphenylmethyl,¹⁷ 1,3-diphenylallyl,^{7,13,18} phenylalkylmethyl,^{14,15} and α - and β -naphthylmethyl¹⁵). However, proton shift changes are quite small and are influenced by factors other than electron density changes, the carbon most sensitive to solvation changes may not be protonated, peak overlap and complex coupling patterns make assignments difficult, and resonances of interest may be obscured by solvent.¹⁷⁻¹⁹ Carbon-13 shifts are a very good empirical measure of π -electron density differences in charged systems.²⁰⁻²⁵ For delocalized carbanions, recent studies have shown that carbon-13 shifts are sensitive to temperature, cation, and solvent.²⁶⁻²⁸

Results and Discussion

Recently, we reported the solvation properties of the alkali metal salts of diphenylmethane using variable-temperature carbon-13 magnetic resonance.²⁹ We report here the completed study of triphenylmethyl (T^-M^+), fluorenyl (F^-M^+),



diphenylmethyl (DPM^-M^+), α -trimethylsilylbenzyl ($TMSB^-M^+$), and benzyl (B^-M^+) alkali metal anions in three solvents of different solvating power: 2-methyltetrahydrofuran (MeTHF), tetrahydrofuran (THF), and dimethoxyethane (DME). Chemical shift changes are very large compared to ¹H NMR shift changes and complex coupling patterns and solvent interference present no problems. Variable temperature carbon-13 magnetic resonance focuses on the changes in electron density that occur as the carbanion undergoes solva-

tion and may be applied to carbanion systems of widely different structure.

General Features of Solvation Changes Observed with Variable-Temperature Carbon-13 Magnetic Resonance. Table I presents carbon-13 chemical shift changes for fluorenyl-lithium dissolved in MeTHF, typical of the changes observed for the other alkali metal carbanions undergoing solvation. These changes are completely reversible within the accessible temperature range. This temperature range is limited at low temperatures by the solubility of the carbanion salt, precipitation usually occurring between -40 and -50 °C. At high temperatures, good spectra are obtained up to just below the solvent's boiling point for all the alkali metal carbanions in MeTHF and THF, and for the lithium salts except benzyllithium in DME. For the sodium, potassium, rubidium, and cesium salts in DME, reaction of the carbanion with solvent becomes rapid above about 60 °C. For benzyllithium in DME, reaction with solvent becomes rapid above about 30 °C.

As the temperature is increased, the α carbon moves upfield and approaches a constant, shielded chemical shift (δ_c), where the equilibrium favors the contact ion. In the contact ion, the cation resides close to the most electron-rich carbon, C- α , and the π -electron density is strongly polarized toward this carbon.³⁰⁻³⁵ In contrast, in the solvent-separated state at low temperatures, the cation is wrapped with solvent and the carbanion responds by spreading the π -electron density more evenly over the carbanion to give a deshielded chemical shift at C- α (δ_s). The largest chemical shift changes occur at the α carbon for all the alkali metal carbanions studied. For fluorenyllithium, the α carbon moves upfield by 5.5 ppm upon an increase in temperature of 130 °C. This change is compensated by a downfield shift change of the two para carbons of 3.5 ppm. Change in π -electron density is the major cause of these shift changes.²⁵ Only minor shift changes are observed at the other carbons (Table I).

Only the lithium and sodium salts undergo appreciable temperature-dependent chemical shift changes. The energy required to convert contact ions into solvent-separated ions is a balance between the coulombic energy needed to separate oppositely charged ions and the energy gained from ion-dipole interaction between the cation and the solvent. The energy needed to separate the cation from the carbanion should decrease as the size of the π system increases and the negative charge becomes more disperse. It is generally accepted that the solvating power of the three solvents is in the order DME \gg THF $>$ MeTHF. Dimethoxyethane has two oxygens fa-

Table I. Variable Temperature Carbon-13 Chemical Shifts for Fluorenyllithium in MeTHF^a

<i>t</i> , °C	C _α	C ₄	C ₁	C ₂	C ₆	C ₃ ^b	C ₅ ^b
-60	82.9	107.8	137.2	116.2	122.5	119.0	118.6
-50	82.8	107.8	137.3	116.2	122.6	119.0	118.6
-40	82.7	107.8	137.3	116.2	122.6	119.0	118.6
-30	82.5	107.8	137.4	116.2	122.6	119.0	118.6
-20	82.2	107.9	137.5	116.2	122.8	119.0	118.7
-10	81.9	108.1	137.6	116.3	122.9	119.1	118.7
0	81.5	108.3	137.6	116.3	123.0	119.3	118.8
10	80.9	108.6	137.5	116.4	123.0	119.4	118.8
20	80.1	109.0	137.4	116.4	123.0	119.7	119.0
30	79.3	109.5	137.3	116.6	123.0	119.9	119.1
40	78.6	109.9	137.0	116.8	122.9	120.2	119.3
50	78.1	110.4	136.8	116.9	122.8	120.3	119.4
60	77.8	110.9	136.5	117.1	122.6	120.7	119.6
70	77.4	111.3	136.3	117.2	122.5	121.0	119.7

^a Chemical shifts are *interpolated* values from measurements taken at 8 °C intervals in parts per million downfield from Me₄Si using cyclohexane as an internal shift reference and the relationship $\delta_{\text{Me}_4\text{Si}} = \delta_{\text{C}_6\text{H}_{12}} + 27.5$. ^b These chemical shift assignments may be interchanged.

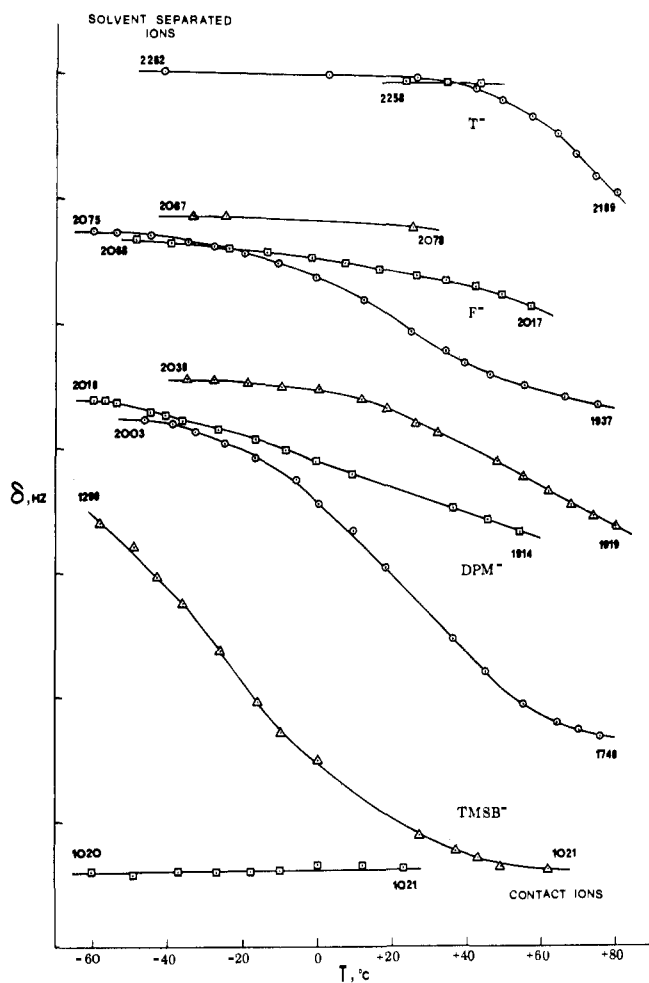


Figure 1. Ion pair equilibrium curves for the lithium salts of T⁻, F⁻, DPM⁻, and TMSB⁻ in MeTHF (○), THF (□), and DME (△). α-Carbon chemical shifts, plotted on the vertical axis, are in hertz downfield from Me₄Si at 25.0335 MHz with 100 Hz between each division. Chemical shifts at the ends of each curve are the α-carbon chemical shifts at the highest and lowest temperatures.

vorably oriented to coordinate electron pairs with the cation. Tetrahydrofuran is superior to MeTHF because the 2-methyl group sterically crowds the coordination site.

The effect of these two features, size of the carbanion's π system and solvating power, is qualitatively demonstrated by the sigmoidal curves for the lithium salts of T⁻, F⁻, DPM⁻,

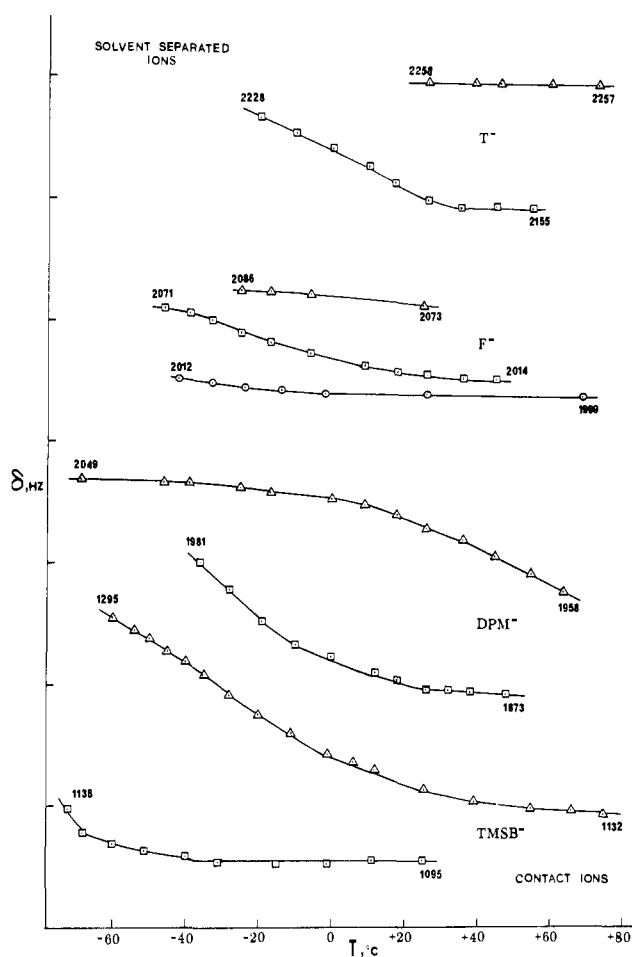


Figure 2. Ion pair equilibrium curves for the sodium salts of T⁻, F⁻, DPM⁻, and TMSB⁻ in MeTHF (○), THF (□), and DME (△). α-Carbon chemical shifts, plotted on the vertical axis, are in hertz downfield from Me₄Si at 25.0335 MHz with 100 Hz between each division. Chemical shifts at the ends of each curve are the α-carbon chemical shifts at the highest and lowest temperatures.

and TMSB⁻ obtained by plotting the α-carbon chemical shifts against temperature (Figure 1). α-Trimethylsilylbenzyl lithium, the smallest π system (8 π electrons, seven carbons), remains a contact ion throughout the temperature range in THF and requires the more powerful solvent, DME, to undergo appreciable solvation. Diphenylmethyl lithium and fluorenyllithium (14 π electrons, 13 carbons) are similar in their

Table II. α -Carbon Chemical Shifts of Contact Ions (δ_c)^a

carbanion	MeTHF δ_c	THF δ_c	DME δ_c	carbanion	MeTHF δ_c	THF δ_c	DME δ_c
T ⁻ Li ⁺	(2706)	(2109)	<i>b</i>	DPM ⁻ Li ⁺	(1682)	(1754)	(1794)
Na ⁺	2129	2155	<i>b</i>	Na ⁺	1830	1873	(1897)
K ⁺	2191	2201	2248	K ⁺	1958	1981	1980
Rb ⁺	2192	2210	<i>c</i>	Rb ⁺	1990	2000	2012
Cs ⁺	2206	2219	<i>c</i>	Cs ⁺	2035	2040	2035
F ⁻ Li ⁺	(1914)	(1946)	<i>c</i>	TMSB ⁻ Li ⁺	1035	1021	1021
Na ⁺	1999	2014	<i>b</i>	Na ⁺	1087	1095	1132
K ⁺	2073	2076	<i>c</i>	K ⁺	1362	1349	1341
Rb ⁺	2094	2084	<i>c</i>	Rb ⁺	1404	1394	<i>c</i>
Cs ⁺	<i>d</i>	2111	<i>c</i>	Cs ⁺	1462	1452	<i>c</i>

^a Chemical shifts in hertz downfield from Me₄Si at 25.0335 MHz; shifts in parentheses estimated from plots of contact shifts of larger alkali metal carbanions vs. 1/interionic distance (Figure 3). ^b Contact shifts not attained at 70 °C and insufficient data for potassium, rubidium, and cesium salts to estimate δ_c . ^c Not determined. ^d Cesium salt insoluble even at high temperatures.

Table III. α -Carbon Chemical Shifts of Lithium and Sodium Solvent-Separated Ions (δ_s) and Estimated Chemical Shift Changes ($\Delta\delta_c$)^a

carbanion	MeTHF		THF		DME	
	δ	$\Delta\delta_c$	δ	$\Delta\delta_c$	δ	$\Delta\delta_c$
T ⁻ Li ⁺	2262	186	2258	149	2260	<i>b</i>
Na ⁺	<i>c</i>		(2258)	103	2258	<i>b</i>
F ⁻ Li ⁺	2075	161	2068	122	2087	<i>b</i>
Na ⁺	<i>c</i>		(2086)	72	2086	<i>b</i>
DPM ⁻ Li ⁺	2003	321	2019	265	2036	242
Na ⁺	<i>c</i>		(2049)	176	2049	152
TMSB ⁻ Li ⁺	<i>c</i>		<i>c</i>		(1452) ^d	431
Na ⁺	<i>c</i>		<i>c</i>		(1452) ^d	320

^a Chemical shifts in hertz downfield from Me₄Si at 25.0335 MHz; shifts in parentheses estimated by comparing to other solvent-separated shifts (see text). ^b Solvent-separated ions; no appreciable change in solvation within accessible temperature range (Figures 1 and 2). ^c Contact ions; no appreciable change in solvation within accessible temperature range (Figures 1 and 2). ^d Complete solvation not attained even at lowest temperatures; shift estimated by using contact ion shift of cesium salt (see text).

solvation characteristics and both are more easily solvated than TMSB⁻Li⁺. In the poorest solvent, MeTHF, they show almost complete change from contact ions to solvent-separated ions; in THF, the middle portions of the solvation curves are observed; and in the most powerful solvent, DME, only the solvent-separated ends of the curves are seen. Triphenylmethyl-lithium, the largest system (20 π electrons, 19 carbons), is the most easily solvated and solvation is complete below about 50 °C except in the poorest solvent, MeTHF.

Figure 2 presents the ion pair equilibrium curves for the sodium salts. Qualitatively, sodium and lithium salts are similar but the lithium salts are slightly easier to solvate. For example, neither TMSB⁻Li⁺ nor TMSB⁻Na⁺ is appreciably solvated in THF and they show about the same extent of solvation in DME. Comparing the larger π systems, it is more apparent that the lithium salts solvate more easily. For example, DPM⁻Li⁺ is completely solvated at low temperatures in MeTHF while DPM⁻Na⁺ requires THF to become only partially solvated. Similarly, T⁻Li⁺ becomes completely solvated in MeTHF while T⁻Na⁺ is only partially solvated in the more powerful solvent, THF.

Carbanions of the larger alkali metals were also studied. Solvation is less thermodynamically favorable than for lithium and sodium cations because the energy gain from the ion-dipole interaction between the more diffuse positive field of the larger cations is much smaller. No significant chemical shift changes were observed for the potassium, rubidium, and cesium salts except for slight downfield changes for DPM⁻K⁺

in THF and DME and for T⁻K⁺ in DME and it must be assumed that the carbanions of the larger alkali metals are contact ions within the accessible temperature range. However, as will be shown, the absence of downfield chemical shift changes at the α carbons may be due to the insensitivity of carbon-13 magnetic resonance to distinguish between contact and solvent-separated ions of the larger alkali metals.

Effect of Cation Size. α -Carbon Chemical Shifts of the Contact Ions (δ_c) and the Solvent-Separated Ions (δ_s). Two very important features emerge when the α -carbon chemical shifts of the contact ions are compared for the same carbanion but for different alkali metal cations. Table II presents observed and estimated δ_c 's for T⁻, F⁻, DPM⁻, and TMSB⁻. First, δ_c is very dependent upon cation size. Second, the size of the chemical shift change experienced by the α carbon upon solvation decreases as the size of the cation increases (Table III, $\Delta\delta_c$). The α -carbon chemical shifts of the contact ions are proportional to the reciprocal of the interionic distance in the contact ion. As the distance between the carbanion and the cation increases, δ_c becomes more deshielded because the positive fields of the larger cations are less effective in polarizing the π -electron density toward the α carbon. Similarly, the absorption maxima of contact ions undergo bathochromic shifts as the size of the cation increases for fluorenyl salts^{1,2,5,36} and alkali metal phenoxides and enolates.^{37,38} However, the change is small and the reciprocal of the cation radius is a sufficient measure of interionic distance to obtain good linear relationships. The UV maxima of ketyls and aromatic hydrocarbon anion radicals are also proportional to the size of the cation in the contact ions.³⁹⁻⁴¹ For these, a distance of 2.0 Å was added to the cation radii to obtain linear plots. Presumably, this 2.0 Å accounts for the contribution of the anion radical to the interionic distance.

The α -carbon chemical shifts of the contact ions also require an estimate of the carbanion's contribution to the interionic distance to be linear and to obtain estimates of the experimentally inaccessible δ_c 's for the lithium salts of T⁻, F⁻, and DPM⁻.⁴² If only the reciprocal of the cation radius is plotted against the observed δ_c 's for Na⁺, K⁺, Rb⁺, and Cs⁺ salts, the least-squares lines predict unreasonably shielded δ_c 's for the lithium salts. For example, the highest field α -carbon chemical shifts for DPM⁻Li⁺ and F⁻Li⁺ in MeTHF at 75 °C are 1749 and 1937 Hz, respectively. Figure 1 clearly shows that these ion pair equilibrium curves have leveled off and only small additional upfield changes are expected. Using only the reciprocal of the cation radius, the least-squares lines predict unreasonably high field values for δ_c : DPM⁻Li⁺, 1537 Hz; F⁻Li⁺, 1837 Hz.

To estimate the interionic distances in the contact ions, the solution radii of these delocalized carbanions are necessary. The best estimates come from conductivity studies and show

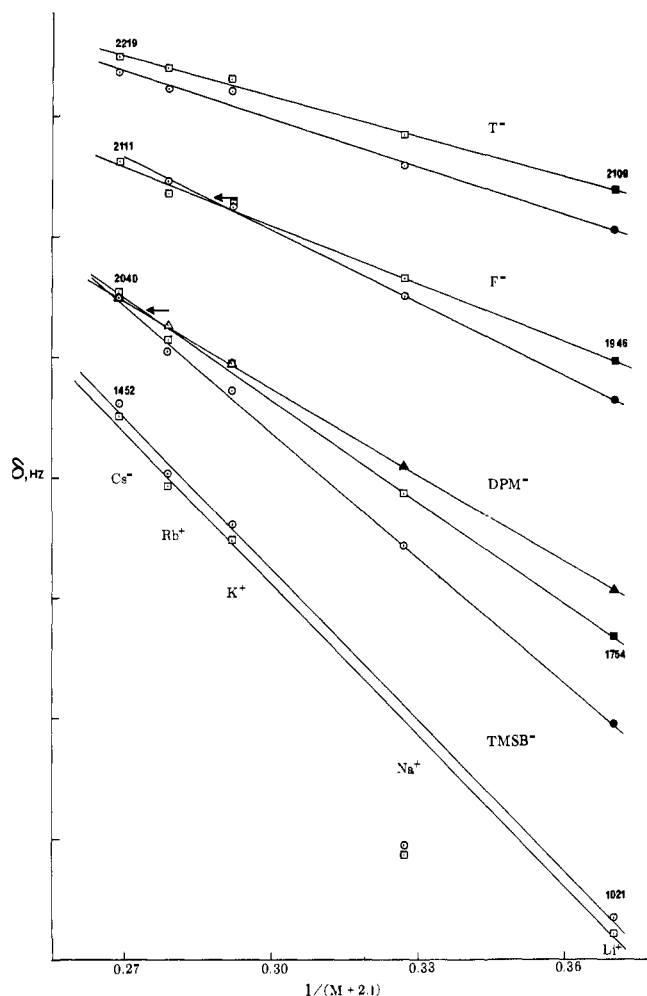


Figure 3. The proportionality between the contact ion α -carbon chemical shifts (δ_c) and interionic distances for T^- , F^- , DPM^- , and $TMSB^-$ in MeTHF (\odot), THF (\square), and DME (Δ). Solid symbols show estimated contact ion shifts. Arrows on the F^- and DPM^- lines indicate approximate position that the horizontal line through the cation-independent lithium and sodium solvent-separated shifts (δ_s) would cross the contact ion lines. For T^- , this shift is downfield from the cesium contact shift at 2260 Hz. Chemical shifts at the ends of each set of curves are the contact shifts in THF.

that the radius of fluorenyl carbanion is 2.1 Å.⁴³ No solution radii have been reported for T^- , DPM^- , and $TMSB^-$. Therefore, we have used 2.1 Å to approximate the radii of all four carbanions. Figure 3 shows the lines that result when the observed δ_c 's are plotted against the reciprocal of the cation radius plus 2.1 Å. Much better predictions are obtained for the inaccessible δ_c 's of the lithium salts. For the lithium salts that are very close to δ_c at high temperatures, the predicted values follow: DPM^-Li^+ , 1682 Hz; F^-Li^+ , 1914 Hz. These estimated contact shift limits were used in the calculations of the thermodynamics of solvation and are believed to be precise to about ± 20 Hz (see Experimental Section). For $TMSB^-Li^+$, all the δ_c 's were experimentally accessible. However, the least-squares lines through the contact shifts for Cs^+ , Rb^+ , K^+ , and Li^+ salts show that the sodium contact shifts are much too shielded. The reason(s) for these deviations is not clear.

Despite the crude method used to estimate the carbanion radii, it must be emphasized that the α -carbon chemical shifts are *proportional* to the interionic distances in the contact ions. The estimates of carbanion radii do not substantially alter the several qualitative conclusions concerning ion-pair structure that may be based upon this important relationship. First, Figure 3 shows the influence of the external solvent shell on the electron distribution in the carbanion. Second, when δ_c 's

of the larger alkali metals are compared to the solvent-separated α -carbon chemical shifts (δ_s 's) of the lithium and sodium salts, this comparison gives a qualitative measure of the change in interionic distance that occurs upon solvation. Third, this relationship provides a warning concerning the use of carbon-13 and proton magnetic resonance for the study of the ion-pair equilibria of carbanions of the larger alkali metals.

Figure 3 shows the influence of solvent structure on the π -electron density at the α carbon for F^- and DPM^- contact ions. The slopes of the lines are shallower for DME than for THF and MeTHF and δ_c is deshielded in the order DME > THF > MeTHF for lithium and sodium cations. DME disperses the cation's charge more effectively than THF or MeTHF because it has two electron donor sites. Therefore, the π -electron density is less polarized toward the α carbon and the δ_c 's are more deshielded in DMF. This small effect of the *external* solvent shell acting through the cation's positive field diminishes as the size of the cation increases and is negligible for K^+ , Rb^+ , and Cs^+ salts. No appreciable solvent effects are seen for the contact ions of either T^- or $TMSB^-$.

As the temperature is lowered, the α -carbon chemical shifts approach a constant, deshielded value (δ_s) as the equilibrium becomes dominated by solvent-separated ions. The solvent-separated α -carbon chemical shifts were accessible for the lithium salts of T^- , F^- , and DPM^- in all solvents and for the sodium salts in DME (Table III, Figures 1 and 2). Unlike the large differences for the δ_c 's, the δ_s 's of the lithium and sodium salts of T^- and F^- in DME are cation *independent* and only a small difference is observed for DPM^- . This implies that solvated lithium and sodium cations are about the same "size". Further, the effect of solvent upon δ_s is negligible for T^- and quite small for F^- and DPM^- (Table III). Therefore, we have estimated the inaccessible δ_s 's for the sodium salts of T^- , F^- , and DPM^- in THF by using the δ_s 's found in DME. These estimates are probably good to about ± 10 Hz. No solvent-separated chemical shifts were observed for $TMSB^-$ even in the most powerful solvent, DME. Therefore, we have used the chemical shift of the contact ion of the cesium salt as a crude estimate of the δ_s 's of the lithium and sodium salts.

Change in Interionic Distance upon Solvation. If the δ_s 's for the smaller lithium and sodium cations are compared to the δ_c 's of the larger alkali metal cations, one may estimate the change in interionic distance that occurs upon solvation. Extending the horizontal line established by the cation-independent lithium and sodium δ_s 's to the contact ion line gives an estimate of the interionic distance for the solvent-separated lithium and sodium ion pairs (Figure 3, \leftarrow). For fluorenyllithium and sodium, the solvent-separated interionic distances are approximately those of fluorenylpotassium contact ion. For diphenylmethyl lithium and sodium, the distances are slightly larger and fall between diphenylmethylrubidium and diphenylmethylcesium contact ions.⁴⁴ In contrast, δ_s 's for triphenylmethyl lithium and sodium (~ 2260 Hz, all solvents) are more deshielded than δ_c of the largest alkali metal contact ion, triphenylmethylcesium (2206, 2219 Hz in MeTHF, THF). This trend ($T^- \gg DPM^- > F^-$) shows the effect of carbanion structure on interionic distances in solvent-separated ion pairs. The solvated cation may approach the bridged, planar fluorenyl anion more closely than diphenylmethyl anion in which rotation about the C- α -C-1 bond makes the anion bulkier. In solvent-separated triphenylmethyl anions, the distance is even larger because of the steric influence of the three phenyl groups. The size of the chemical shift change at the α carbon depends upon the size of the π system and the change in interionic distance upon solvation (Table III, $\Delta\delta_c: DPM^- > T^- > F^-$). For the 14- π -electron systems, DPM^- undergoes a larger change than F^- because of the larger interionic distance required in the solvent-separated ion. For the 20- π -electron system, T^- , one might expect the smallest $\Delta\delta$ because the

Table IV. Equilibrium Constants (SS/CI) and Thermodynamics of Solvation^a

solvent carbanion	equilibrium constants				Δt	$\Delta\delta_0$	%	ΔH , kcal/mol	ΔS , deg · cal/ mol	ref
	-25°C	0°C	25°C	50°C						
MeTHF										
T ⁻ Li ⁺	SS	SS	SS	7.41	31	93	50	-15.5 ± 0.8	-44 ± 3	<i>b</i>
F ⁻ Li ⁺	>10	3.35	1.01	0.36	66	138	86	-7.7 ± 1.0	-26 ± 4	<i>c</i>
DPM ⁻ Li ⁺	>10	3.79	1.31	0.76	93	254	79	-7.1 ± 0.5	-23 ± 2	
THF										
T ⁻ Na ⁺	2.89	0.84	<0.10	CI	50	73	70	-11.7 ± 0.7	-44 ± 2	<i>d</i>
F ⁻ Li ⁺	>10	6.17	3.36	1.84	50	51	42	-4.9 ± 0.3	-14 ± 2	<i>e</i>
F ⁻ Na ⁺	1.18	0.30	<0.10	CI	57	56	78	-7.9 ± 0.9	-32 ± 3	<i>f</i>
DPM ⁻ Na ⁺	0.80	0.21	<0.10	CI	54	108	61	-7.6 ± 0.3	-31 ± 1	
DME										
DPM ⁻ Li ⁺	SS	>10	6.33	2.58	62	117	48	-7.2 ± 0.2	-20 ± 1	
DPM ⁻ Na ⁺	>10	7.59	3.00	1.19	81	91	60	-6.6 ± 0.4	-20 ± 2	
TMSB ⁻ Li ⁺	0.67	0.26	<0.10	CI	95	278	64	-5.3 ± 0.3	-23 ± 1	
TMSB ⁻ Na ⁺	0.42	0.19	<0.10	CI	72	163	43	-3.5 ± 0.2	-16 ± 1	

^a Equilibrium constants are *interpolated* values obtained from ion pair equilibrium curves (Figures 1 and 2); $\Delta\delta_0$ is the observed α -carbon chemical shift change used for van't Hoff plots over the temperature range, Δt ; % is the percent conversion within the temperature range, Δt ($\Delta\delta_0/\Delta\delta_c$). ^b Lit.^{5c} -11.8 kcal/mol and -32 cal/deg-mol in cyclohexylamine. ^c Lit.^{5c} -8.5 kcal/mol and -23 cal/deg-mol. ^d Lit.¹⁷ -8.2 kcal/mol and -28 cal/deg-mol. ^e Lit.^{1,2b} -7.5 kcal/mol and -27 cal/deg-mol. ^f Lit.¹⁷ -6.7 kcal/mol and -27 cal/deg-mol; lit.^{1,2b} -7.6 kcal/mol and -33 cal/deg-mol.

larger π system should be less susceptible to polarization by the cation or the solvated cation. However, this is compensated for by the larger interionic distance required in the solvent-separated ion.

Solvation of the Larger Alkali Metal Cations. For the same carbanion, $\Delta\delta$ decreases greatly from lithium to sodium. From Figure 3, one may estimate that if a potassium salt were to undergo solvation, the chemical shift change would be quite small (<30 Hz). This observation provides a warning concerning the use of magnetic resonance measurements to study the solvation properties of carbanions of the larger alkali metal cations.⁴⁵ The trends established by the lithium and sodium salts and simple thermodynamic arguments allow the conclusion that solvation is much less favorable because the energy gain from ion-dipole interaction between the large cation and solvent is small. Therefore, it is safe to assume that the K⁺, Rb⁺, and Cs⁺ salts are contact ions. However, it is important to recognize that if solvation were to occur for these cations, the change in chemical shift at the most sensitive α carbon might not be observable.

Ion Pair Equilibrium Constants and Thermodynamics of Solvation. For carbanions that undergo at least 40% of the conversion from solvent-separated ions to contact ions, van't Hoff plots of the equilibrium constants vs. the reciprocal of the absolute temperature give good straight lines (Table IV, $K = (\delta_0 - \delta_c)/(\delta_s - \delta_0)$ where δ_0 = observed α -carbon chemical shift, δ_c = α -carbon contact ion chemical shift, and δ_s = α -carbon solvent-separated ion chemical shift). These thermodynamic values are in close agreement with those of T⁻ and F⁻ salts determined by UV^{1,2b,5} and ¹H NMR¹⁷ measurements except for F⁻Li⁺ in THF. For F⁻Li⁺ in THF, the values are probably low because the solvent-separated limit must be estimated, the chemical shift change is relatively small, and only 42% of the conversion from contact to solvent-separated ions can be observed. The enthalpies of solvation depended upon cation size and are similar for lithium and sodium. With the largest π system, T⁻, and the smallest, TMSB⁻, the enthalpy's dependence upon carbanion structure is apparent. The exceptionally large enthalpies for trityl salts are due to a combination of two factors. The energy necessary for cation separation is smaller because of the large π system and the energy gain from cation solvation is greater because of the large interionic distance. The salts of the smallest carbanion, TMSB⁻Li⁺ and TMSB⁻Na⁺, have enthalpies lower than for DPM⁻ and F⁻. This is due to the more concentrated, less

polarizable π system which increases the energy debt for cation separation.

The large, negative entropies reflect the increased order upon solvation with the "freezing out" of solvent molecules around the cation. These entropies show a solvent dependence similar to that previously reported for fluorenyl salts in THF and DME.¹⁷ Since DME has two coordination sites, fewer molecules of solvent are immobilized than in THF or MeTHF and the entropies are usually smaller in DME. The large entropies for the trityl salts may be associated with the large interionic distances required in the solvent-separated ions.

Solvation of Benzyl Carbanions. Variable temperature carbon-13 spectra were obtained for benzyl lithium, sodium, potassium, and rubidium. The alkali metal carbanions of α -trimethylsilyltoluene were surprisingly stable toward protonation by solvent compared to the benzyl salts. Good spectra were obtained for TMSB⁻M⁺ at high temperatures even in DME, the solvent most sensitive to deprotonation. In contrast, benzyl carbanions could be studied only in MeTHF and THF at ambient temperatures and below. Spectra were obtained for benzyl lithium in DME up to 30 °C if the carbanion was prepared at low temperatures. The greater stability of TMSB⁻ in ethereal solvents is undoubtedly due to the delocalizing influence of the trimethylsilyl group attached to the electron-rich α carbon.⁴⁶

Contact ions dominate at high temperatures as the size of the π system decreases (Figures 1 and 2). TMSB⁻Li⁺ and TMSB⁻Na⁺ are primarily contact ions above about 30 °C in all solvents. Therefore, it is safe to assume that the benzyl anions are contact ions at room temperature. Further, it can be predicted that benzyl lithium and benzyl sodium would undergo very large chemical shift changes upon solvation. An estimate of this change is obtained by comparing the chemical shift of the contact ion of benzyl lithium to that of benzyl rubidium and assuming that solvated lithium will have about the same "size" as a rubidium cation: 54.6 - 36.5 = 18.1 ppm, 453 Hz. The variable temperature α -carbon chemical shifts for benzyl anions are presented in Table V. For benzyl lithium in THF and DME, only very small downfield chemical shift changes are observed at low temperatures. Therefore, it may be concluded that very little solvation occurs for benzyl lithium or the other benzyl alkali metal anions. For this small carbanion system, the energy of cation separation is comparable to the energy gain upon solvation and solvation is much less thermodynamically favorable.

Table V. α -Carbon Chemical Shifts for Benzyl Carbanions^a

solvent <i>t</i> , °C	THF Li ⁺ δ	DME Li ⁺ δ	25 °C	THF δ
25	915	906	Na ⁺	976
0	921	<i>b</i>	K ⁺	1319
-25	927	<i>b</i>	Rb ⁺	1366
-50	936	982		
$\Delta\delta_0$	21	76		

^a Chemical shifts in hertz downfield from Me₄Si at 25.0335 MHz.

^b Not determined.

Experimental Section

Materials. Solvents were refluxed over potassium metal for 24 h, distilled under argon, and then maintained at reflux over sodium-potassium alloy in a solvent still. Lithium sand was available as a 20% dispersion in mineral oil. Sodium and potassium sand were prepared under mineral oil after the removal of the surface oxide coat by heating above the melting points and vigorous shaking as the sample cooled. Mineral oil was removed by several washings with dry solvent. Rubidium and cesium metal were stored in sealed ampules and the ampules broken in the carbanion sample tubes under dry cyclohexane. Hydrocarbon precursors were dried with sodium sand prior to use. Dibenzylmercury was prepared by the reaction of benzylmagnesium chloride with mercuric chloride and recrystallized from ethanol, and the last traces of ethanol were removed under vacuum.

Bis(trimethylsilylmethyl)mercury ([Me₃SiCH₂)₂Hg). A mixture of bromomethyltrimethylchlorosilane and chloromethyltrimethylchlorosilane was prepared by the addition of bromine to an excess of trimethylchlorosilane while chlorine was bubbled through the solution and the reaction flask was irradiated with a sunlamp.⁴⁷ The reaction mixture was fractionated and the fraction between 95 and 134 °C (740 mm) collected to give a 49% yield based upon bromine. The mixed halosilanes were reacted with methylmagnesium bromide to give a mixture of halomethyltrimethylsilanes boiling between 95 and 120 °C, 80% yield [lit.⁴⁸ bp bromo, 115 °C (742 mm); bp chloro, 97 °C (734 mm)]. Bis(trimethylsilylmethyl)mercury was prepared by the addition of mercuric chloride dissolved in THF to the Grignard reagent of the mixed halomethyltrimethylsilanes.⁴⁹ Fractional distillation gave bis(trimethylsilylmethyl)mercury, 89% yield, bp 58–62 °C (2 mm) (lit.⁵⁰ bp 49–50 °C).

Bis(benzyltrimethylsilylmethyl)mercury ([PhCH₂SiMe₂CH₂)₂Hg). Halomethyltrimethylbenzylsilane was prepared by the reaction of benzylmagnesium chloride with the mixed halomethyltrimethylchlorosilanes. The Grignard reagent of halomethyltrimethylbenzylsilane was added to mercuric chloride to give bis(benzyltrimethylsilylmethyl)mercury, 20% yield, mp 49–51 °C. ¹H NMR (integrated intensity): -SiMe₂, δ -0.02 (6); (-CH₂)₂Hg, -0.28 (2); PhCH₂-, 2.03 (2); aryl, 7.13 (5). ¹³C NMR (multiplicity): -SiMe₂, δ 1.02 (4, ³J_{C-Hg}¹⁹⁹ = 46 Hz); (-CH₂)₂Hg, 23.8 (3, ¹J_{C-Hg}¹⁹⁹ = 580 Hz); PhCH₂-, 28.6 (3); aromatic carbons, C-1, 140.7 (1); C-2, 127.7 (2); C-3, 128.1 (2); C-4, 123.7 (2).

Carbanion Preparations. Carbanion solutions were prepared in septum-capped, argon-flushed, 12-mL centrifuge tubes. All transfers of solvent and carbanion solutions were made with oven-dried glass syringes. The concentration of the carbanion solutions was varied between 0.2 and 0.7 M. Within this range, no significant chemical shift changes were noted for the same temperature for different concentrations of the same carbanion. Three methods were used to prepare the carbanions: metal reduction of hydrocarbon precursors; metal reduction of organomercurials; and metal reduction of bis(trimethylsilylmethyl)mercury followed by metalation of hydrocarbon precursors.

Direct metal reduction was used to prepare F⁻Li⁺, F⁻Na⁺, and F⁻K⁺. The centrifuge tube containing alkali metal sand (~12 mg-atoms), solvent (8 mL), and fluorene (1–5 mmol) was placed in an ice bath in an ultrasonic bath for approximately 1 h. The ultrasonic bath was particularly valuable in initiating these heterogeneous reductions. Good spectra were not obtained for F⁻Rb⁺ and F⁻Cs⁺ by metal reduction of fluorene, probably owing to the formation of paramagnetic species due to further reduction.

All the alkali metal carbanions of triphenylmethane and diphenylmethane and the rubidium and cesium salts of fluorene were pre-

pared by the metalation of the hydrocarbon precursors with trimethylsilylmethylmetal compounds directly in MeTHF, THF, or DME. Bis(trimethylsilylmethyl)mercury (0.5–2.5 mmol) was syringed into the centrifuge tube containing alkali metal sand (~12 mg-atoms), solvent (8 mL), and hydrocarbon precursor (1–5 mmol). Reduction of the organomercurial was initiated at 0 °C in an ultrasonic bath and the carbanions were completely formed in approximately 0.5 h. The potassium, rubidium, and cesium salts of benzyltrimethylsilyl, benzylpotassium, and benzylrubidium were prepared in a similar manner except that the reduction of the organomercurial and the metalation of the precursor were carried out heterogeneously in cyclohexane, cyclohexane was removed, and ethereal solvent was added slowly at -78 °C.

The lithium and sodium salts of benzyltrimethylsilyl were prepared by metal sand reduction of bis(benzyltrimethylsilylmethyl)mercury at 0 °C. In MeTHF, the initially formed benzyltrimethylsilyllithium (PhCH₂SiMe₂CH₂⁻Li⁺) could be observed in the carbon-13 spectrum and it was necessary to heat the sample to 50 °C for 1 h to complete the rearrangement to α -trimethylsilylbenzyl lithium (PhCH⁻SiMe₂, Li⁺). In THF and DME, this rearrangement was rapid and only the α -trimethylsilylbenzyl anion was observed spectroscopically. Benzyl lithium and benzyl sodium were prepared by the metals and reduction of dibenzylmercury at -30 °C.

Carbon-13 Measurements. All carbon-13 spectra were obtained with a JEOL PFT-100/Nicolet 1080 Fourier transform spectrometer with an internal D₂O capillary for locking. At the same temperature, chemical shifts were reproducible to about ± 4 Hz. Approximately 1000–3000 average transients were required to obtain satisfactory signal to noise ratios for proton-decoupled spectra. Chemical shifts were referenced to tetramethylsilane and to cyclohexane (689.1 Hz at 25.0335 MHz from Me₄Si). It is important not to use solvent resonances as references, particularly the carbon attached directly to the solvent's oxygen. These resonances varied by as much as ± 20 Hz and depended upon carbanion, cation, and temperature. Temperature of the carbanion sample was controlled to ± 1 °C and was measured before and after each spectral measurement.

Error Estimates in the Calculation of Thermodynamic Values. Calculations of thermodynamic values and equilibrium constants were performed for lithium and sodium salts that covered at least 40% of the change from one solvation state to the other and where one of the chemical shift limits was reached at high or low temperature. Calculations were not carried out for the larger alkali metal salts because only very small chemical shift changes were observed (e.g., DPM⁻K⁺, THF, $\Delta\delta$ = 13 Hz, ΔT = 65 °C; DPM⁻K⁺, DME, 24 Hz, 80 °C; T⁻K⁺, DME, 17 Hz, 74 °C) and the precision of such calculations is likely to be poor because the range covered if solvation occurs would be small.

The largest contributing factor to the error estimates comes from the estimates of the inaccessible chemical shift limits. The contact ion chemical shift limits were estimated from least-squares plots of the contact limits of the larger alkali metal salts vs. $1/(M + 2.1 \text{ \AA})$ (Figure 3). The errors presented in Table IV for salts with estimated δ_c 's were obtained by varying δ_c by ± 20 Hz (T⁻Li⁺, MeTHF, 2076 \pm 20; F⁻Li⁺, MeTHF, 1914 \pm 20; F⁻Li⁺, THF, 1946 \pm 20; DPM⁻Li⁺, MeTHF, 1682 \pm 20; DPM⁻Li⁺, DME, 1794 \pm 20; DPM⁻Na⁺, DME 1897 \pm 20). Solvent-separated chemical shift limits are less subject to error because they are cation independent. Therefore, for carbanions that required estimated δ_s 's, errors were obtained by varying the estimated limit by ± 10 Hz (T⁻Na⁺, THF, 2258 \pm 10; F⁻Na⁺, THF, 2086 \pm 10; DPM⁻Na⁺, THF, 2049 \pm 10). For TMSB⁻, the δ_s 's were crudely estimated because no solvent-separated limits were accessible. Therefore, error estimates were obtained by varying the estimated limits by ± 40 Hz (TMSB⁻Li⁺, DME, 1452 \pm 40; TMSB⁻Na⁺, DME, 1452 \pm 40). It should be noted that, even though the solvent-separated limits for TMSB⁻ are crude, the errors are comparable to those for the other carbanions because the error caused by estimating the solvent-separated limit is largely compensated for by the very large change covered by TMSB⁻ upon solvation.

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- (45) Since the chemical shift changes are much smaller, this warning is likely to be even more important for the use of 1H NMR in solvation studies (ref 17 and 19).
- (46) From the carbon-13 spectra of nine different benzyl anions compared to the carbon-13 spectra of benzyl anions with a trimethylsilyl group attached to the electron-rich α carbon, we find that the π systems in the trimethylsilyl-substituted carbanions are deshielded by 20.8 ± 0.8 ppm. Since a trimethylsilyl deshields carbon by about 6 ppm when it replaces hydrogen, approximately 15 ppm reflects silicon's ability to delocalize electron density through $d\pi-p\pi$ bonding. Silicon-29 shifts are shielded by 16-26 ppm in the silyl anions compared to their precursors.
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